Since the actually measured inductance, $L_{2^{\circ}act}$, of the coil is a function that can be expressed also with respect to the absolute magnetic permeability of 'the vacuum', as $L_{2^{\circ}act} = \mu_0 N^2 A_2/h$, where by AToS, $\mu_0 = W_x^{-2}$, the function i_{MB} can also be expressed in the form of a proportional relation between the two area functions, A_1 and A_2 :

$$\begin{split} \mathbf{i}_{\text{MB}} &= \mathbf{N} \ \mathbf{B}_{2^{\circ}\text{MB}} \ \mathbf{A}_{1}/\mathbf{L}_{2^{\circ}} = 4\pi^{2} \ [(\mathbf{B}_{2^{\circ}\text{MB}} \ \mathbf{h})/(\boldsymbol{\mu}_{o} \ \mathbf{N})] \ \mathbf{A}_{1}/\mathbf{A}_{2} = \\ &= (4\pi^{2} \ \mathbf{A}_{1}/\mathbf{N}) \ \mathbf{B}_{2^{\circ}\text{MB}} \ \mathbf{h} \ \mathbf{W}_{x}^{2}]/\mathbf{A}_{2} = [(4\pi^{2} \ \mathbf{A}_{1}/\mathbf{N})/(\mathbf{A}_{2}] \ \mathbf{B}_{2^{\circ}\text{MB}} \ \mathbf{h} \ \alpha^{-1} \ 10^{2} \ c^{2} = \\ &= [(10^{*} \ 2^{*} \ \pi)^{2} \ \alpha^{-1}]^{*} \ [\mathbf{B}_{2^{\circ}\text{MB}} \ \mathbf{h} \ c^{2} \ (\mathbf{A}_{1}/\mathbf{A}_{2} \ast \mathbf{N})] \end{split}$$

At this juncture, we can present a schematic comparison between classical electromagnetic theory of the operation of induction coils, and the AToS approach to what is the synchronous induction of two distinct charge fluxes, one massbound and the other massfree, and thus propose a general function for alternating nonelectromagnetic currents, where all charge is deployed by longitudinal electric wave functions and associated transverse magnetic wave functions:



From the above, it becomes evident how the classical relations purport to account for the